

MIDTERM I

(A) PICK STANDARD PRIMITIVE VECTORS

$$\vec{a}_1 = L \left(\frac{\sqrt{3}}{2} \hat{x} + \frac{1}{2} \hat{y} \right)$$

$$\vec{a}_2 = L \left(\frac{\sqrt{3}}{2} \hat{x} - \frac{1}{2} \hat{y} \right)$$

WITH A 2-ATOM BASIS $\{0, a\hat{x}\}$

NOW NEED TO FIGURE OUT HOW L IS RELATED TO A.

FROM GEOMETRY, ENY TO SEE

$$\frac{L}{2} = a \cos 30^\circ = a \frac{\sqrt{3}}{2} \Rightarrow L = \sqrt{3}a$$

SO

$$\vec{a}_1 = a \left(\frac{3}{2} \hat{x} + \frac{\sqrt{3}}{2} \hat{y} \right); \quad \vec{a}_2 = a \left(\frac{3}{2} \hat{x} - \frac{\sqrt{3}}{2} \hat{y} \right)$$

BASIS $\{0, a\hat{x}\}$

(B) (SEE SHEET)

(C) HONEYCOMB CRYSTAL IS TRIANGULAR (HEXAGONAL) ^{LATTICE} WITH BASIS. HENCE, RECIPROCAL LATTICE IS TRIANGULAR (HEXAGONAL). (THERE REALLY IS NO HONEYCOMB LATTICE)

TO FIND THE RECIPROCAL LATTICE VECTORS,

JUST ADD A THIRD ONE,

$\vec{a}_3 = c\hat{z}$ AND USE FORMULA

FROM A&M FOR 3-D, BUT

KEEP ONLY THE ONES IN PLANE.

$$\Rightarrow \vec{b}_1 = 2\pi \frac{\vec{a}_2 \times \vec{a}_3}{\vec{a}_1 \cdot \vec{a}_2 \times \vec{a}_3}$$

$$\vec{b}_2 = 2\pi \frac{\vec{a}_3 \times \vec{a}_1}{\vec{a}_1 \cdot \vec{a}_2 \times \vec{a}_3}$$

FROM PART (A)

$$\vec{a}_2 \times \vec{a}_3 = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{3}{2}a & -\frac{\sqrt{3}}{2}a & 0 \\ 0 & 0 & c \end{vmatrix} = -\frac{\sqrt{3}}{2}ac\hat{x} - \frac{3}{2}ac\hat{y}$$

$$\vec{a}_3 \times \vec{a}_1 = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 0 & 0 & c \\ \frac{3}{2}a & \frac{\sqrt{3}}{2}a & 0 \end{vmatrix} = \frac{\sqrt{3}}{2}ac\hat{x} + \frac{3}{2}ac\hat{y}$$

AND

$$\left(\vec{a}_1 \cdot \vec{a}_2 \times \vec{a}_3 \right) = \begin{vmatrix} \frac{3}{2}a & \frac{\sqrt{3}}{2}a & 0 \\ \frac{3}{2}a & -\frac{\sqrt{3}}{2}a & 0 \\ 0 & 0 & c \end{vmatrix} = -\frac{3\sqrt{3}}{2} a^2 c$$

So

$$\vec{b}_1 = \frac{2\pi}{a} \left(\frac{1}{3} \hat{x} + \frac{\sqrt{3}}{3} \hat{y} \right)$$

$$\vec{b}_2 = \frac{2\pi}{a} \left(\frac{1}{3} \hat{x} - \frac{\sqrt{3}}{3} \hat{y} \right)$$

$$\textcircled{D} \quad S_{\vec{k}} = \sum_{\substack{\{\text{BASIS}\} \\ \vec{d}}} e^{i\vec{k} \cdot \vec{d}}$$

$$\vec{d} = \xi \vec{0}, \quad a \hat{x} \xi$$

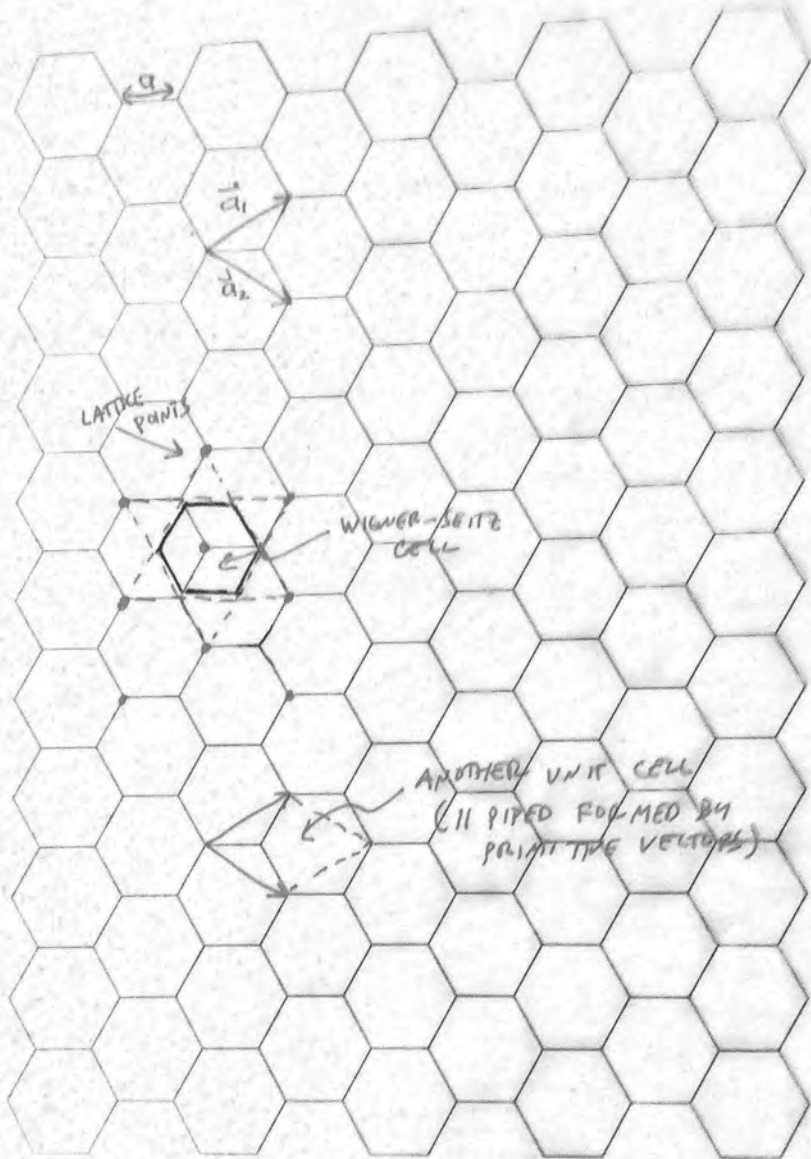
$$\vec{k} = n_1 \vec{b}_1 + n_2 \vec{b}_2$$

So

$$S_{\vec{k}} = 1 + \text{EXP} \left(i (n_1 + n_2) \frac{2\pi}{3} \right)$$

So,

$$S_{\vec{k}} = \begin{cases} 2 & (n_1 + n_2) \bmod 3 = 0 \\ 1 + e^{i\frac{2\pi}{3}} & (n_1 + n_2) \bmod 3 = 1 \\ 1 + e^{i\frac{4\pi}{3}} & (n_1 + n_2) \bmod 3 = 2 \end{cases}$$



DEPENDENCE OF VOLUME OF PRIMITIVE VECTORS

②

FOR BCC, a_{bcc} IS THE LENGTH OF THE CUBE (NOT NEAREST NEIGHBOR DISTANCE)

$$V_{bcc} = a_{bcc}^3 \quad \text{WITH } N=2 \text{ ATOMS PER UNIT CELL}$$

FOR HCP, THERE ARE 2 ATOMS PER UNIT CELL (ONE AT

LATTICE POINT AND ONE AT $\frac{a}{\sqrt{3}}\hat{x} + \frac{c}{2}\hat{z}$). FOR HCP, $\frac{c}{a} = \frac{\sqrt{8}}{\sqrt{3}}$

PRIMITIVE VECTORS hcp!

$$\vec{a}_1 = a_{hcp} \hat{x}$$

$$\vec{a}_2 = a_{hcp} \left(\frac{\hat{x}}{2} + \frac{\sqrt{3}\hat{y}}{2} \right)$$

$$\vec{a}_3 = c \hat{z} = \sqrt{8/3} a_{hcp} \hat{z}$$

VOLUME UNIT CELL IS

$$\vec{a}_1 \cdot \vec{a}_2 \times \vec{a}_3 = \begin{vmatrix} a_{hcp} & 0 & 0 \\ a_{hcp}/2 & \frac{\sqrt{3}}{2} a_{hcp} & 0 \\ 0 & 0 & \sqrt{8/3} a_{hcp} \end{vmatrix} = a_{hcp}^3 \sqrt{2}$$

SINCE $N_{bcc} = N_{hcp} = 2$

AND DENSITY IS CONSTANT,

$\Rightarrow V_{hcp} = V_{bcc} \Rightarrow$

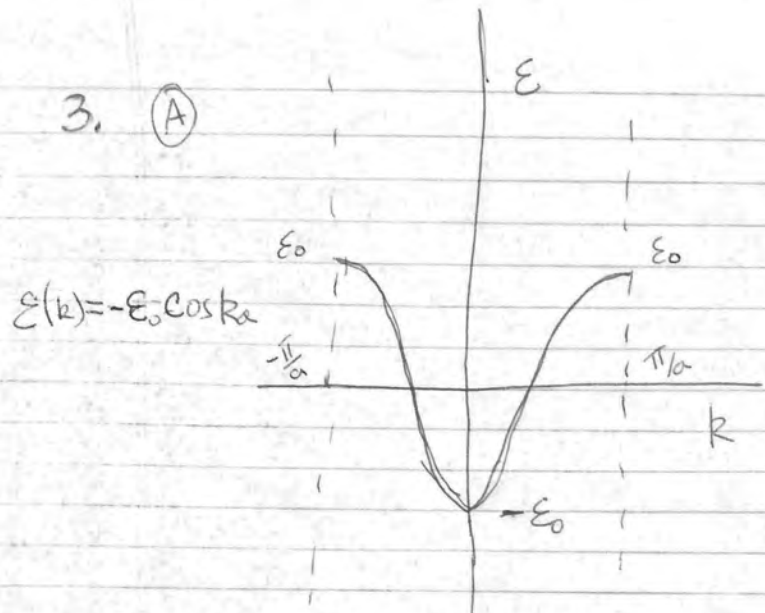
$$a_{hcp}^3 \sqrt{2} = a_{bcc}^3$$

OR $a_{hcp} = \frac{1}{\sqrt[3]{2}} a_{bcc}$

FOR $a_{bcc} = 4.23 \text{ \AA} \Rightarrow$

$$a_{hcp} = 3.77 \text{ \AA}$$

3. (A)



(B)

$$\frac{1}{m^*} = \frac{1}{\hbar^2} \frac{\partial^2 E}{\partial k^2}$$

$$\frac{\partial^2 E}{\partial k^2} = E_0 a^2 \cos ka$$

$$\text{SO } m^* = \left(\frac{\hbar^2}{E_0 a^2} \right) (\cos ka)^{-1}$$

@ $k=0$, $m^* = \frac{\hbar^2}{E_0 a^2}$

AT BZ BOUNDARY ($k = \pm \frac{\pi}{a}$)

$$m^* = -\frac{\hbar^2}{E_0 a^2}$$

(C) PERIODIC BOUNDARY CONDITIONS \Rightarrow

$$\psi_k(0) = \psi_k(L)$$

$$\Rightarrow e^{ikL} = 1 \Rightarrow$$

$$kL = m(2\pi)$$

AND $k = \frac{m2\pi}{L}$ WITH m AN INTEGER

THE NUMBER OF "k" STATES BETWEEN

$$k = -\frac{\pi}{a} \text{ AND } k = +\frac{\pi}{a}$$

TOLLS US m CAN GO FROM

$$m = -\frac{L}{2a} \dots +\frac{L}{2a}$$

SO THERE ARE $2\left(\frac{L}{2a}\right) = \left(\frac{L}{a}\right)$

ALLOWED STATES, IF WE INCLUDE SPIN, WE GET

(YOU CAN COUNT THE $m=0$ STATE TOO, BUT REMEMBER $k=+\pi/a$ AND $k=-\pi/a$ ARE THE SAME STATE SINCE THEY DIFFER BY $2\pi/a$.)

$$N = \frac{2L}{a} \text{ STATES (WITH SPIN)}$$

FOR THE "LINEAR" DENSITY,

$$n = \frac{N}{L} = \frac{2}{a}$$

(D) WE HAVE AT $T=0$

$$N = \sum_{k,s} 1 = 2 \left(\frac{L}{2\pi}\right) \int_{-k_F}^{+k_F} dk$$

\uparrow SPIN $\underbrace{\hspace{2cm}}$ DENSITY OF STATES IN k -SPACE

HENCE

$$N = 4\left(\frac{L}{2\pi}\right) k_F \text{ SO}$$

$$k_F = \left(\frac{N}{L} \frac{\pi}{2}\right) = n \frac{\pi}{2}$$

$$k_F = n \frac{\pi}{2} \quad | - D$$

$$\text{So } \boxed{\epsilon_F = \epsilon(k_F) = -\epsilon_0 \cos\left(\frac{n\pi a}{2}\right)}$$

(n MUST BE BELOW $n_{\text{MAX}} = \frac{2}{a}$
TO BE IN THE LOWEST BAND)

$$\textcircled{E} \quad \sum_{\text{SPIN}} \sum_k \rightarrow Z \left(\frac{L}{2\pi}\right) \int_{-k}^k dk \rightarrow$$

$$Z \left(\frac{L}{2\pi}\right) \int_0^k dk \rightarrow \int_0^{\epsilon} D(\epsilon) d\epsilon$$

↑
INTEGRATE
OVER $+k$ ONLY

$$\text{So } \underbrace{D(\epsilon) d\epsilon}_{\text{DENSITY OF STATES IN ENERGY}} = Z \left(\frac{L}{2\pi}\right) dk$$

$$\underbrace{g(\epsilon)}_{\text{DENSITY OF STATES/UNIT LENGTH}} = \frac{D(\epsilon)}{L} = \frac{Z}{\pi} \frac{dk}{d\epsilon}$$

$$g(\epsilon) = \frac{Z}{\pi} \frac{1}{\left| \frac{d\epsilon}{dk} \right|} \quad \frac{d\epsilon}{dk} = \epsilon_0 a \sin ka$$

$$\text{So } \boxed{g(\epsilon) = \frac{Z}{\pi \epsilon_0 a |\sin ka|}}$$

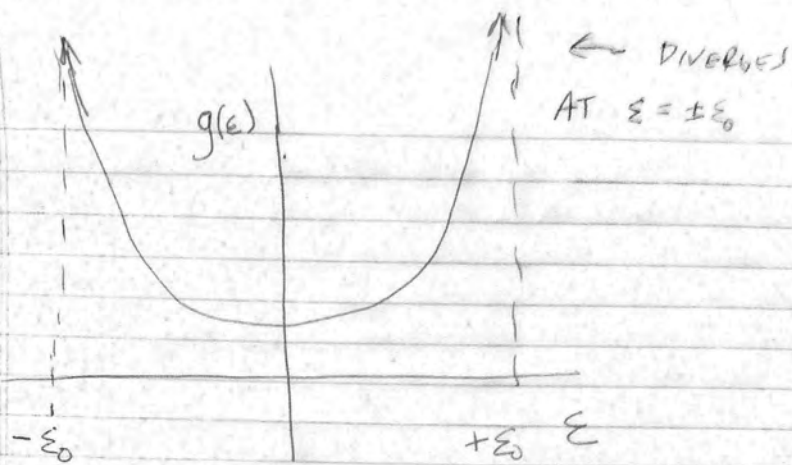
NOTE! ① HAS RIGHT UNITS
ENERGY-LENGTH

② WE STILL HAVE k IN THE EXPRESSION (I WILL ACCEPT THIS).

REALLY, WE SHOULD ELIMINATE k IN TERMS OF ϵ

$$|\epsilon_0 \sin ka| = \sqrt{\epsilon_0^2 - \epsilon^2}$$

$$\boxed{g(\epsilon) = \frac{Z}{\pi a} \frac{1}{\sqrt{\epsilon_0^2 - \epsilon^2}}}$$



4. THE RECIPROCAL LATTICE VECTORS FOR A BCC CRYSTAL ARE:

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2.31

$$\vec{b}_1 = \frac{2\pi}{a} (\hat{y} + \hat{z}) \quad \vec{b}_2 = \frac{2\pi}{a} (\hat{x} + \hat{z})$$

$$\vec{b}_3 = \frac{2\pi}{a} (\hat{x} + \hat{y})$$

SO A GENERAL RECIPROCAL LATTICE VECTOR IS:

n_1, n_2, n_3
INTEGERS

$$\vec{G} = n_1 \vec{b}_1 + n_2 \vec{b}_2 + n_3 \vec{b}_3$$

$$\vec{G} = \frac{2\pi}{a} \left((n_2 + n_3) \hat{x} + (n_1 + n_3) \hat{y} + (n_1 + n_2) \hat{z} \right)$$

$$n_1, n_2, n_3 \in \mathbb{Z}$$

THE ENERGY BANDS IN THE DIRECTION FROM Γ TO H HAVE A WAVE VECTOR

$$\vec{k} = \frac{2\pi}{a} (k \hat{x}) \quad \text{WITH } k=0 \text{ TO } 1$$

THE EMPTY LATTICE ENERGY BANDS ARE

$$E = \frac{\hbar^2}{2m} (\vec{k} + \vec{G})^2$$

WE CAN SCALE THE ENERGY AND WAVEVECTOR TO GET

$$\tilde{E} = \frac{E}{\frac{\hbar^2}{2m} \left(\frac{2\pi}{a}\right)^2} = \left((k_x + n_1 + n_2 + n_3)^2 + (n_1 + n_3)^2 + (n_1 + n_2)^2 \right)$$

WITH k GOING FROM 0 TO 1 AND $n_1, n_2, n_3 \in \mathbb{Z}$

THE ENERGY THAT IS LOWEST AT THE "H" POINT IS $n_1 = n_2 = n_3 = 0$ AND CORRESPONDS TO $\tilde{E} = 1$ (scaled)

WE THEN NEED TO GO TO $\tilde{E} = 5$

WE MAKE A TABLE LIKE THE ONE ON PAGE 176.

** NOTE: (n_1, n_2, n_3) ARE THE INTEGERS TIMES $\frac{a}{2\pi}$, $\frac{a}{2\pi}$, $\frac{a}{2\pi}$

NOT $\hat{x}, \hat{y}, \hat{z}$

$\tilde{E} = \frac{E}{\frac{\hbar^2}{2m} \left(\frac{2\pi}{a}\right)^2}$	BAND	$G \frac{a}{2\pi}$	$\tilde{E}(0,0,0)$	$\tilde{E}(k_x, 0, 0)$
1		(0,0,0)	0	\tilde{k}_x^2
2		(1,0,0)	2	$2 + \tilde{k}_x^2$
3		(1,0,0)	2	
4		(0,1,0)	2	$(k_x + 1)^2 + 1$
5		(0,0,1)	2	"
6		(0,1,0)	2	$(k_x - 1)^2 + 1$
7		(0,0,1)	2	"
8		(0,1,1)	2	$2 + \tilde{k}_x^2$
9		(0,1,1)	2	

	10	$(1, \bar{1}, 0)$	2	} $(\tilde{k}_x - 1)^2 + 1$
	11	$(1, 0, \bar{1})$	2	
	12	$(0, 1, 1)$	6	} $(\tilde{k}_x + 2)^2 + 2$
	13	$(0, \bar{1}, \bar{1})$		
≥ 5		$(1, 1, 0)$	6	} $(\tilde{k}_x + 1)^2 + 5$
		$(1, 0, 1)$		
≥ 5		$(\bar{1}, \bar{1}, 0)$		
		$(\bar{1}, 0, \bar{1})$		} $(\tilde{k}_x - 1)^2 + 5$
	14	$(\bar{1}, 1, 0)$	2	
	15	$(\bar{1}, 0, 1)$	2	} $(\tilde{k}_x + 1)^2 + 1$
≥ 5		$(1, 1, 1)$	12	} $(\tilde{k}_x \pm 2)^2 + 8$
		$(\bar{1}, \bar{1}, \bar{1})$		
	16	$(1, \bar{1}, 1)$	4	$\tilde{k}_x^2 + 4$
	17	$(1, 1, \bar{1})$		
	18	$(\bar{1}, 1, 1)$	4	$(\tilde{k}_x \pm 2)^2$
	19	$(1, \bar{1}, \bar{1})$	4	

20	$(\bar{1}, 1, \bar{1})$	4	$\tilde{k}_x^2 + 4$
21	$(\bar{1}, \bar{1}, 1)$		

OUCH!! NOW - THE LESS, THERE ARE ONLY A FEW PLOTS WE NEED TO DO!

BOUND	\tilde{k}_x^2	
1	\tilde{k}_x	
2, 3, 8, 9	$2 + \tilde{k}_x^2$	
4, 5, 14, 15	$(\tilde{k}_x + 1)^2 + 1$	
6, 7, 10, 11	$(\tilde{k}_x - 1)^2 + 1$	
12	$(\tilde{k}_x + 2)^2 + 2$	} NOTE THESE ARE < 5 FOR $\tilde{k}_x \neq 0$
13	$(\tilde{k}_x - 2)^2 + 2$	
16, 17, 20, 21	$\tilde{k}_x^2 + 4$	
18	$(\tilde{k}_x + 2)^2$	
19	$(\tilde{k}_x - 2)^2$	

$$\frac{\hbar^2}{2m} \left(\frac{\pi}{a} \right)^2$$

