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▷ Conductors.

↳ Materials in which charges can move freely.

- electrons in metallic conductors

- ions in ionic conductor such as salt water.

• an ideal conductor or perfect conductor has 0 resistivity.

$\rho \rightarrow 0$ resistivity $\rightarrow \sigma \rightarrow \infty$ ($\rho = \frac{1}{\sigma}$)
 \uparrow conductivity

Here, our discussion is on an ideal conductor.

(i) $\vec{E} = 0$ inside a conductor: No force on charges

(ii) No excess charges inside a conductor: $\rho = 0$
 $\vec{\nabla} \cdot \vec{E} = \rho / \epsilon_0$ ($\vec{E} = 0$ inside) \rightarrow

(iii) Any net charge resides on the surface.

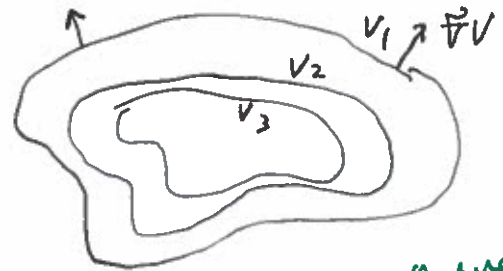
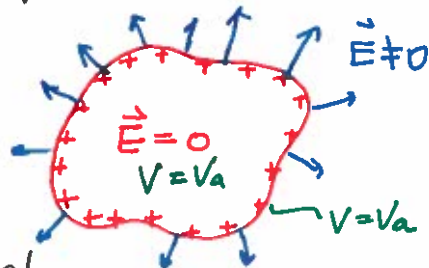
(iv) A conductor is an equipotential

$$V(b) - V(a) = - \int_a^b \vec{E} \cdot d\vec{r} = 0 \text{ inside a conductor.}$$

\Rightarrow Surface of a conductor is an equipotential surface

(v) Therefore, \vec{E} on the surface is perpendicular to the surface.

$$-\vec{\nabla} V(r) = \vec{E}$$

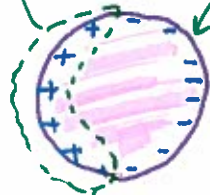


▷ Induced Charges.

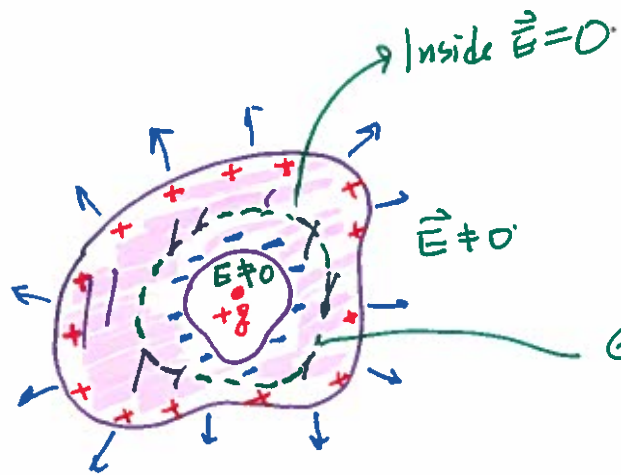
deficiency of (-) charge from neutral \Rightarrow (+) effective.
 induced (-) charge attracted by +q



neutral conductor no net charge



+q

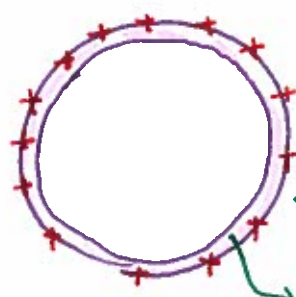


Gaussian surface
 Inside $q_{en} = +q + q_{ind} = 0$
 $q_{ind} = -q$

In a perfect conductor, the induced charge completely cancel out the presence of $+q$ in the cavity

EX

(1)

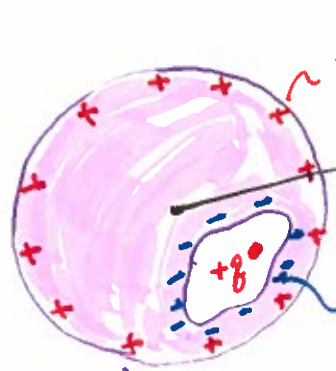


total net charge on a conductivity shell
 $= +q$

only on the outer surface.

$\vec{E} = 0$ inside the shell. (no charge on the inner surface)

(2)



$$\vec{E} = \frac{1}{4\pi\epsilon_0} \cdot \frac{+q}{r^2}$$

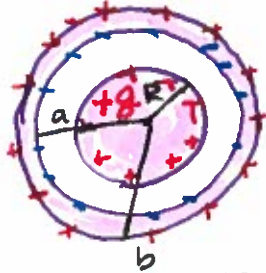
$-q$ (completely screen $+q$ out)

neutral conductivity sphere w/ cavity.

HW

2.38 & 2.39

2.38



$$\sigma_R = \frac{+Q}{4\pi R^2}; \quad \sigma_a = \frac{-Q}{4\pi a^2}; \quad \sigma_b = \frac{+Q}{4\pi b^2}$$

$$\vec{E} = E \hat{r} \quad \text{where} \quad E(r) = \begin{cases} 0 & (r < a) \\ \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} & (a < r < b) \\ 0 & (a < r < b) \\ \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} & (r > b) \end{cases}$$

$$V(0) = - \int_{\infty}^{n=0} \vec{E} \cdot d\vec{r} = - \int_{\infty}^b \frac{Q}{4\pi\epsilon_0} \cdot \frac{1}{r^2} dr + \int_b^a 0 dr + \int_a^R \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} dr + \int_b^R 0 dr$$

$$= \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{b} + \frac{1}{R} - \frac{1}{a} \right)$$

if you connect the outer shell to GND, then the charge (+) will be drained to GND. $\sigma_b \rightarrow 0$, consequently $E(r) = 0$ for $r > a$.

$$\therefore V(0) = \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{R} - \frac{1}{a} \right)$$

► Surface charge and the force on a conductor.

Boundary conditions.

$$\vec{E}_{\text{out}} - \vec{E}_{\text{in}} = \frac{\sigma}{\epsilon_0} \hat{n} \quad \Rightarrow \text{"Discontinuity in } \vec{E} \text{ at the surface charge"}$$

outside conductor *inside conductor* $\vec{E}_{\text{in}} = 0$.

$$\vec{E}_{\text{out}} = \vec{E} = \frac{\sigma}{\epsilon_0} \hat{n}$$

$$\text{or } \frac{\partial V_{\text{out}}}{\partial n} - \frac{\partial V_{\text{in}}}{\partial n} = -\frac{1}{\epsilon_0} \sigma \quad \left(\frac{\partial V}{\partial n} = \vec{\nabla} V \cdot \hat{n} \right)$$

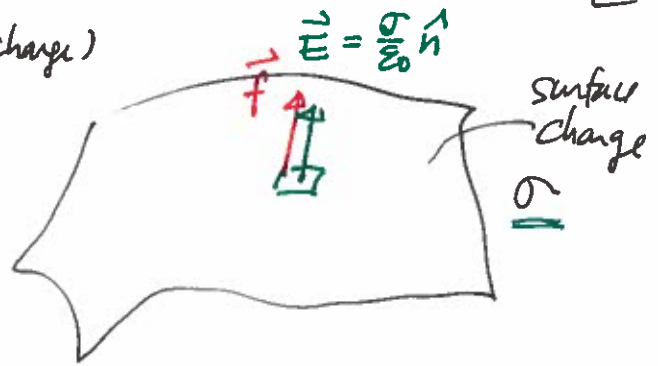
$\underbrace{\frac{\partial V_{\text{in}}}{\partial n}}_{=0 \text{ (equipotential)}}$ \approx normal derivative.

$$\therefore \frac{\partial V_{\text{out}}}{\partial n} - \frac{\partial V}{\partial n} = -\frac{1}{\epsilon_0} \sigma$$

presence of \vec{E} where σ (surface charge)

\Rightarrow force on σ

\vec{f} : force on the surface charge of unit area. $\therefore \vec{f} = \sigma(\vec{v}) = \sigma$.



$\therefore \vec{f} = \sigma \vec{E} \rightarrow$ what is the proper \vec{E} hence?

$$\vec{E} = \vec{E}_{avg} = \frac{1}{2} (\vec{E}_{out} + \vec{E}_{in}) = \frac{1}{2} \frac{\sigma}{\epsilon_0} \hat{n}$$

$$\therefore \vec{f} = \frac{1}{2} \frac{\sigma^2}{\epsilon_0} \hat{n}$$

$$\therefore d\vec{f} = \frac{1}{2} \frac{\sigma^2}{\epsilon_0} da \hat{n} = \frac{1}{2} \frac{\sigma^2}{\epsilon_0} d\vec{a}$$

For an arbitrary surface of a conductor, the total force

$$\vec{F} = \int_S d\vec{f} = \frac{1}{2} \int_S \frac{\sigma^2}{\epsilon_0} d\vec{a} \quad (\sigma \text{ could be non-uniform.})$$

► Capacitor (more than one conductor)

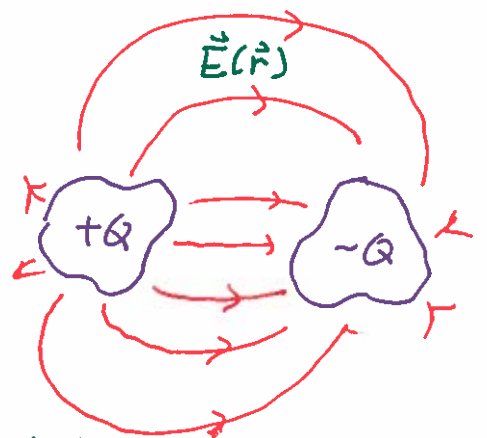
The electric pot. difference b/w two conductors

$$\Delta V = V_+ - V_- = - \int_{(-)}^{(+)} \vec{E} \cdot d\vec{r}$$

Depends on the exact distribution of surface charges on the conductors

\Downarrow geometrical effect!

But we know $|\vec{E}| \propto Q$.



$\therefore \Delta V \propto Q \cdot \left[\text{geometrically determined factor} \right]$

$\boxed{\square} \equiv \frac{1}{C}$ where C is called capacitance.

$\therefore \Delta V = \frac{Q}{C}$ or $Q = C \Delta V$.

for a given ΔV , the amount of charge on the conductor is ~~directly~~ proportional to C (capacitance).

$[C] = \left[\frac{Q}{V} \right] = \frac{C}{V} = F$ (farad).

Q What is the sign of ΔV ?
Always $\Delta V > 0$ \therefore

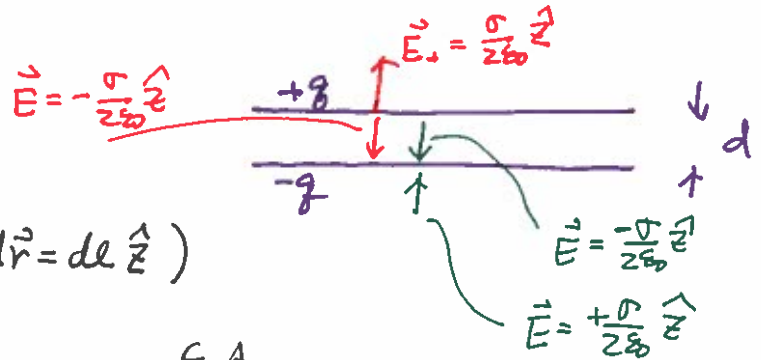
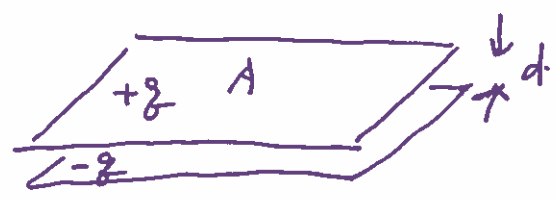
Q How to calculate a capacitance?

- (i) Assign $+q$ & $-q$ on each conductor.
- (ii) Calculate \vec{E} and $\Delta V = V_+ - V_-$.
- (iii) Then you'll establish $\Delta V = \boxed{\square} q$
 \downarrow
 $\frac{1}{C}$

EX (1) Parallel plate capacitor

Between the plates

$\vec{E} = -\frac{\sigma}{\epsilon_0} \hat{z}$ ($\sigma = \frac{q}{A}$)
 $= -\frac{q}{\epsilon_0 A} \hat{z}$



$\therefore \Delta V = V_+ - V_-$

$= -\int_{(-)}^{(+)} \vec{E} \cdot d\vec{r}$ ($d\vec{r} = dl \hat{z}$)

$= \frac{d}{\epsilon_0 A} \cdot q$

$\therefore C = \frac{\epsilon_0 A}{d}$

(2) Concentric Spherical Shells.

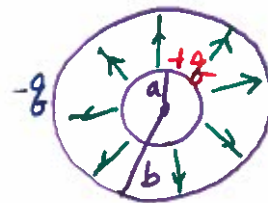
$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r} \quad (\text{between the shell}).$$

$$\therefore \Delta V = V_+ - V_- = - \int_{(c)}^{(+)} \frac{q}{4\pi\epsilon_0} \frac{1}{r^2} \hat{r} \cdot d\vec{r}$$

$$= \frac{-q}{4\pi\epsilon_0} \int_b^a \frac{1}{r^2} dr \quad (d\vec{r} = dr \hat{r})$$

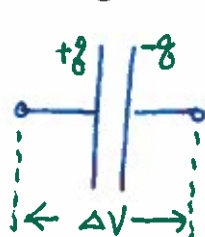
$$= \frac{q}{4\pi\epsilon_0} \left(\frac{1}{r} \right) \Big|_b^a = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{a} - \frac{1}{b} \right)$$

$$\therefore C = 4\pi\epsilon_0 \cdot \frac{ab}{b-a} \quad (> 0).$$



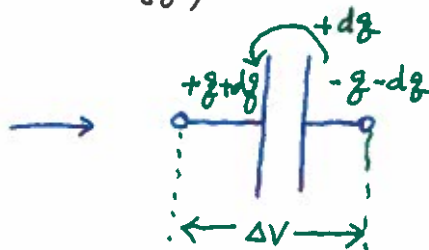
HW 2.43

⊙ Energy (electrostatic energy) stored in a capacitor.



Capacitance: C

$$\Delta V = \frac{q}{C}$$



The amount of work to move $+dq$ from (-) plate to (+) plate for an arbitrary charge q

$$dW = \Delta V dq = \frac{q}{C} dq$$

\therefore total amount of work to establish charge q from 0 (nothing)

$$W = \int_0^q \frac{q'}{C} dq' = \frac{1}{2} \frac{q^2}{C} \quad \text{or}$$

$$= \frac{1}{2} C (\Delta V)^2$$

HW 2.44 and 2.54 (additional)

$$[W] = [E] \quad [Force] = [Pressure] \cdot [Area] \quad \therefore [pressure] = \frac{[Energy]}{[Volume]}$$