

L4

Ch. 2.

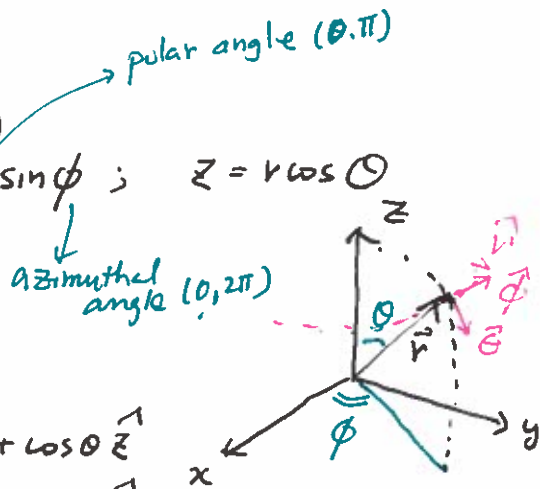
► Curvilinear Coordinates

⊙ Spherical Coordinates (r, θ, φ)

- $x = r \sin \theta \cos \varphi$; $y = r \sin \theta \sin \varphi$; $z = r \cos \theta$

$$\vec{A} = A_r \hat{r} + A_\theta \hat{\theta} + A_\varphi \hat{\varphi}$$

$$= A_r \hat{e}_r + A_\theta \hat{e}_\theta + A_\varphi \hat{e}_\varphi$$



$$\hat{r} = \sin \theta \cos \varphi \hat{x} + \sin \theta \sin \varphi \hat{y} + \cos \theta \hat{z}$$

$$\hat{\theta} = \cos \theta \cos \varphi \hat{x} + \cos \theta \sin \varphi \hat{y} - \sin \theta \hat{z}$$

$$\hat{\varphi} = -\sin \varphi \hat{x} + \cos \varphi \hat{y}$$

$$\hat{r} \cdot \hat{\theta} = \hat{\theta} \cdot \hat{\varphi} = \hat{r} \cdot \hat{\varphi} = 0$$

$$\hat{r} \times \hat{\theta} = \hat{\varphi} ; \hat{\theta} \times \hat{\varphi} = \hat{r} ; \hat{\varphi} \times \hat{r} = \hat{\theta}$$

"orthogonal coordinate system"

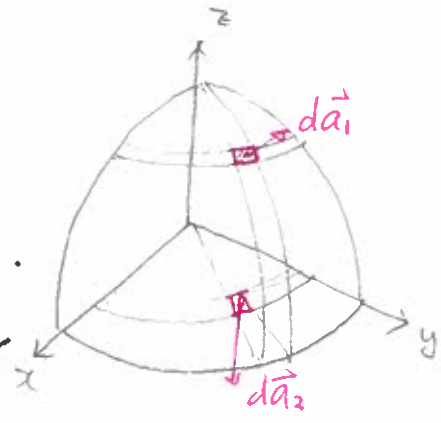
* all function of θ, φ .
not fixed in space!

$$d\vec{r} = d\vec{e} = dr \hat{r} + r d\theta \hat{\theta} + r \sin \theta d\varphi \hat{\varphi}$$

$$dV = d\tau = r^2 \sin \theta dr d\theta d\varphi$$

$$d\vec{a}_1 = r^2 \sin \theta d\theta d\varphi \hat{r}$$

$$d\vec{a}_2 = r dr d\varphi \hat{\theta}$$



EX Volume of a sphere by integration.

$$V = \int d\tau = \int_0^R \int_{\theta=0}^{\pi} \int_{\varphi=0}^{2\pi} r^2 \sin \theta d\theta d\varphi dr$$

$$= \int_0^R r^2 dr \int_0^{\pi} \sin \theta d\theta \int_0^{2\pi} d\varphi$$

$$= \frac{1}{3} R^3 \cdot 2 \cdot \pi = \frac{4}{3} \pi R^3$$

$$\vec{\nabla} F = \left(\frac{\partial F}{\partial x}\right) \hat{x} + \left(\frac{\partial F}{\partial y}\right) \hat{y} + \left(\frac{\partial F}{\partial z}\right) \hat{z}$$

$$= \square \hat{r} + \triangle \hat{\theta} + \circ \hat{\varphi}$$

$$x = x(r, \theta, \varphi); \quad y = y(r, \theta, \varphi); \quad z = z(r, \theta, \varphi).$$

$$\frac{\partial}{\partial x} = \frac{\partial}{\partial r} \frac{\partial r}{\partial x} + \frac{\partial}{\partial \theta} \frac{\partial \theta}{\partial x} + \frac{\partial}{\partial \varphi} \frac{\partial \varphi}{\partial x}$$

$$\frac{\partial}{\partial y} = \dots$$

$$\frac{\partial}{\partial z} = \dots$$

$$\hat{x} = \odot \hat{r} + \square \hat{\theta} + \triangle \hat{\varphi}$$

$$\hat{y} = \dots$$

$$\hat{z} = \dots$$

$\frac{\partial r}{\partial r} = 1$	$\frac{\partial \theta}{\partial r} = 0$	$\frac{\partial \varphi}{\partial r} = 0$
$\frac{\partial r}{\partial \theta} = 0$	$\frac{\partial \theta}{\partial \theta} = 1$	$\frac{\partial \varphi}{\partial \theta} = 0$
$\frac{\partial r}{\partial \varphi} = 0$	$\frac{\partial \theta}{\partial \varphi} = 0$	$\frac{\partial \varphi}{\partial \varphi} = 1$
$\frac{\partial r}{\partial \varphi} = \sin \theta \hat{\varphi}$	$\frac{\partial \theta}{\partial \varphi} = \cos \theta \hat{\varphi}$	
	$\frac{\partial \varphi}{\partial \varphi} = -\sin \theta \hat{r} - \cos \theta \hat{\theta}$	

Finally,

$$\vec{\nabla} F = \frac{\partial F}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial F}{\partial \theta} \hat{\theta} + \frac{1}{r \sin \theta} \frac{\partial F}{\partial \varphi} \hat{\varphi}$$

$$\vec{\nabla} \cdot \vec{A} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta A_\theta) + \frac{1}{r \sin \theta} \frac{\partial A_\varphi}{\partial \varphi}$$

$$\vec{\nabla} \times \vec{A} = \frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (\sin \theta A_\varphi) - \frac{\partial A_\theta}{\partial \varphi} \right] \hat{r} + \frac{1}{r} \left[\frac{1}{\sin \theta} \frac{\partial A_r}{\partial \varphi} - \frac{\partial}{\partial r} (r A_\varphi) \right] \hat{\theta}$$

$$+ \frac{1}{r} \left[\frac{\partial}{\partial r} (r A_\theta) - \frac{\partial A_r}{\partial \theta} \right] \hat{\varphi}$$

$$\nabla^2 F = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial F}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial F}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 F}{\partial \varphi^2}$$

⊙ Cylindrical coordinates (s, φ, z)

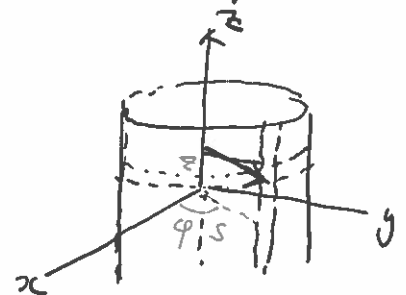
$$x = s \cos \varphi; \quad y = s \sin \varphi; \quad z = z.$$

$$\vec{A} = A_s \hat{s} + A_\varphi \hat{\varphi} + A_z \hat{z} \quad (\hat{s} \perp \hat{\varphi} \perp \hat{z})$$

$$\hat{s} = \cos \varphi \hat{x} + \sin \varphi \hat{y}$$

$$\hat{\varphi} = -\sin \varphi \hat{x} + \cos \varphi \hat{y}$$

$$\hat{z} = \hat{z}$$



$$d\vec{r} = d\vec{\ell} = ds \hat{s} + s d\varphi \hat{\varphi} + dz \hat{z}$$

$$dV = d\tau = s ds d\varphi dz$$

$$\vec{\nabla} F = \left(\frac{\partial F}{\partial s}\right) \hat{s} + \frac{1}{s} \left(\frac{\partial F}{\partial \varphi}\right) \hat{\varphi} + \left(\frac{\partial F}{\partial z}\right) \hat{z}$$

$$\vec{\nabla} \cdot \vec{A} = \frac{1}{s} \frac{\partial}{\partial s} (s A_s) + \frac{1}{s} \frac{\partial A_\varphi}{\partial \varphi} + \frac{\partial A_z}{\partial z}$$

$$\vec{\nabla} \times \vec{A} = \left(\frac{1}{s} \frac{\partial A_z}{\partial \varphi} - \frac{\partial A_\varphi}{\partial z}\right) \hat{s} + \left(\frac{\partial A_s}{\partial z} - \frac{\partial A_z}{\partial s}\right) \hat{\varphi} + \frac{1}{s} \left[\frac{\partial}{\partial s} (s A_\varphi) - \frac{\partial A_s}{\partial \varphi}\right] \hat{z}$$

$$\nabla^2 F = \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial F}{\partial s}\right) + \frac{1}{s^2} \frac{\partial^2 F}{\partial \varphi^2} + \frac{\partial^2 F}{\partial z^2}$$

o Generalization of Coordinate Systems ('orthogonal')

- Suppose that every point in 3D space can be represented by three coordinates (u_1, u_2, u_3) .
 u_i 's are smooth functions of $\vec{r} = (x, y, z)$.

- The essence of any coordinate system is in the differential length interval, ds .

$$ds^2 = (h_1 du_1)^2 + (h_2 du_2)^2 + (h_3 du_3)^2$$

where h_i are non-negative scalar factor, varying smoothly with \vec{r} or u_i

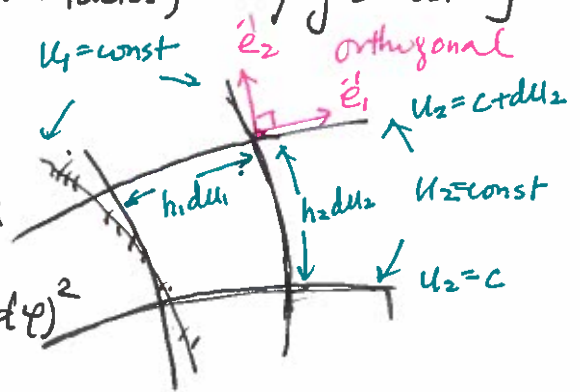
For example, in spherical coord.

~~$$d\vec{\ell} = (dr \hat{r} + r d\theta \hat{\theta} + r \sin \theta d\varphi \hat{\varphi})$$~~

$$d\vec{\ell} = dr \hat{r} + r d\theta \hat{\theta} + r \sin \theta d\varphi \hat{\varphi}$$

$$\therefore (ds)^2 = d\vec{\ell} \cdot d\vec{\ell} = (dr)^2 + (r d\theta)^2 + (r \sin \theta d\varphi)^2$$

$$\therefore h_r = 1 ; h_\theta = r ; h_\varphi = r \sin \theta$$



example of 2D curvilinear system

Also in spherical cylindrical coord.

$$h_s = 1 ; h_\varphi = s ; h_z = 1$$

- All vector derivatives can be written the following general forms.

$$\vec{\nabla} F = \frac{1}{h_1} \frac{\partial F}{\partial u_1} \hat{e}_1 + \frac{1}{h_2} \frac{\partial F}{\partial u_2} \hat{e}_2 + \frac{1}{h_3} \frac{\partial F}{\partial u_3} \hat{e}_3$$

$$\vec{\nabla} \cdot \vec{A} = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} (h_2 h_3 A_1) + \dots \right]$$

$$\vec{\nabla} \times \vec{A} = \frac{1}{h_2 h_3} \left[\frac{\partial}{\partial u_2} (h_3 A_3) - \frac{\partial}{\partial u_3} (h_2 A_2) \right] \hat{e}_1 + \dots$$

$$\nabla^2 F = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} \left(\frac{h_2 h_3}{h_1} \frac{\partial F}{\partial u_1} \right) + \dots \right]$$

HW Prob. 1.38 & 1.39

$$\vec{r} = x \hat{x} + y \hat{y} + z \hat{z}$$

$$= r \sin \theta \cos \varphi \hat{x} + r \sin \theta \sin \varphi \hat{y} + r \cos \theta \hat{z}$$

$$\hat{r} = \frac{d\vec{r}}{dr} / \left| \frac{d\vec{r}}{dr} \right| ; \hat{\theta} = \frac{d\vec{r}}{d\theta} / \left| \frac{d\vec{r}}{d\theta} \right| ; \hat{\varphi} = \frac{d\vec{r}}{d\varphi} / \left| \frac{d\vec{r}}{d\varphi} \right|$$

meaning of the expression.

HW Prob. 1.43

HW Prob. 1.62

▷ Dirac Delta function

○ 1D Dirac delta fun.

$$\delta(x) = \begin{cases} 0 & \text{if } x \neq 0 \\ \infty & \text{if } x = 0 \end{cases}$$

with $\int_{-\infty}^{\infty} \delta(x) dx = 1$

- $\delta(-x) = \delta(x)$: symmetric fun.

- If $f(x)$ is continuous,

$$f(x)\delta(x) = f(0)\delta(x)$$

$$\therefore \int_{-\infty}^{+\infty} f(x)\delta(x)dx = f(0) \int_{-\infty}^{+\infty} \delta(x)dx = f(0)$$

And also

$$\int_{-\infty}^{+\infty} f(x)\delta(x-a)dx = f(a) \int_{-\infty}^{+\infty} \delta(x-a)dx = f(a)$$

$$- \delta(kx) = \frac{1}{|k|}\delta(x)$$

$$- x \frac{d}{dx}(\delta(x)) = -\delta(x)$$

HW Prob. 1.44

o 3D Dirac delta function .

$$\delta^3(\vec{r}) = \delta(x)\delta(y)\delta(z) = \begin{cases} 0 & \text{otherwise} \\ \infty & \text{if } x=y=z=0 \end{cases}$$

$$\text{with } \int_{-\infty}^{+\infty} dx dy dz \delta^3(\vec{r}) = 1$$

In general for d -dim

$$\delta^d(\vec{r}) = \prod_{i=1}^d \delta(x_i)$$

$$- \int_V f(\vec{r})\delta^3(\vec{r}-\vec{a})dV = f(\vec{a})$$

$$- \vec{\nabla} \cdot \left(\frac{\hat{r}}{r^2} \right) = 4\pi\delta^3(\vec{r}) \quad \text{and} \quad \vec{\nabla} \cdot \left(\frac{\hat{r}}{r^2} \right) = 4\pi\delta^3(\vec{r})$$

$$\text{We know } \vec{\nabla} \cdot \left(\frac{\hat{r}}{r} \right) = -\frac{\hat{r}}{r^2}$$

$$\therefore \vec{\nabla} \cdot \vec{\nabla} \left(\frac{1}{r} \right) = \nabla^2 \left(\frac{1}{r} \right) = -\vec{\nabla} \cdot \left(\frac{\hat{r}}{r^2} \right) = -4\pi\delta^3(\vec{r}) \quad **$$

$$- \delta^3(k\vec{r}) = \frac{1}{|k|^3}\delta^3(\vec{r})$$

* Early HV problem (1.16). $\vec{\nabla} \cdot \left(\frac{\hat{r}}{r^2} \right) = \vec{\nabla} \cdot \left(\frac{\vec{r}}{r^3} \right)$

In Cartesian coordinate

$$\vec{\nabla} \cdot \left(\frac{\hat{r}}{r^2} \right) = \frac{\partial}{\partial x} \left(\frac{x}{(x^2+y^2+z^2)^{3/2}} \right) + \dots$$

But in spherical coordinate

$$\vec{\nabla} \cdot \left(\frac{\hat{r}}{r^2} \right) = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \cdot \frac{1}{r^2} \right) = 0 \quad \text{for } r \neq 0.$$

$\frac{\hat{r}}{r^2}$ is not well-defined at $r=0$.

$$\int_V \vec{\nabla} \cdot \left(\frac{\hat{r}}{r^2} \right) d\tau = \oint_S \frac{\hat{r}}{r^2} \cdot d\vec{s} \quad \left(d\vec{s} \text{ on the spherical surface of radius } r \Rightarrow R^2 \sin\theta d\theta d\phi \hat{r} \right)$$

$$= \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi \frac{\hat{r} \cdot \hat{r}}{R^2} R^2 = 4\pi.$$

$$\int_V 4\pi \delta^3(\vec{r}) d\tau = 4\pi.$$

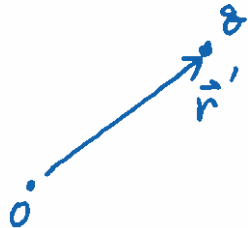
EX

$$J = \int_V (r^2+z) \vec{\nabla} \cdot \left(\frac{\hat{r}}{r^2} \right) d\tau$$

$$= \int_V (r^2+z) 4\pi \delta^3(\vec{r}) d\tau = 8\pi.$$

EX Consider the following point charge configurations.

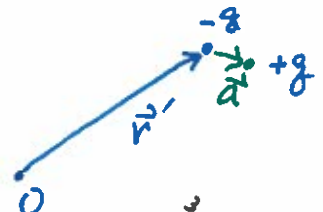
(i) Point charge at \vec{r}'



charge density $\rho(\vec{r}) = q \delta^3(\vec{r} - \vec{r}')$

$$\int_V \rho(\vec{r}) d\tau = \int q \delta^3(\vec{r} - \vec{r}') d\tau = q$$

(ii) Dipole at \vec{r}' separated by \vec{a}



$\rho(\vec{r}) = -q \delta^3(\vec{r} - \vec{r}') + q \delta^3(\vec{r} - \vec{r}' - \vec{a})$

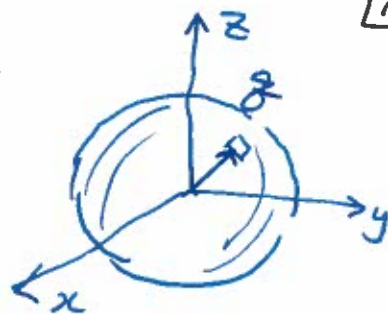
$$\int_V \rho(\vec{r}) d\tau = -q + q = 0$$

(iii) ^{Uniform} Surface charge on a sphere of radius R .

$$\rho(\vec{r}) = \frac{q}{4\pi R^2} \delta(r-R)$$

1D delta fun.

$$\int_V \rho(\vec{r}) d\tau = 4\pi R^2 \int \frac{q}{4\pi R^2} \delta(r-R) dr = q$$



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HW Prob. 1.48 (b) & (c).

► Delta function & Fourier transformation.

⊙ Let $f(x)$ be a fun. of a real variable x , and as $x \rightarrow \pm\infty$ $|x|^n f(x) \rightarrow 0$ (~~is bounded~~) for any $n > 0$. Then we can define the Fourier transform of $f(x)$: f_k

$$\left. \begin{aligned} f_k &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x) e^{-ikx} dx \text{ and also inversely} \\ f(x) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f_k e^{+ikx} dx \end{aligned} \right\} \text{--- ①}$$

For a fun in 3D, $f(\vec{r})$, one can define the F.T.

$$f_{\vec{k}} = \frac{1}{(2\pi)^{3/2}} \int d^3r f(\vec{r}) e^{-i\vec{k}\cdot\vec{r}} \text{ and } f(\vec{r}) = \frac{1}{(2\pi)^{3/2}} \int d^3k f_{\vec{k}} e^{+i\vec{k}\cdot\vec{r}}$$

From Eq. ①

$$\begin{aligned} f_k &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x) e^{-ikx} dx = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} dx e^{-ikx} \cdot \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} dk' f_{k'} e^{+ik'x} \\ &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} dk' f_{k'} \int_{-\infty}^{+\infty} dx e^{+i(k'-k)x} \\ &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} dk' f_{k'} \underbrace{2\pi \delta(k'-k)} = f_k. \end{aligned}$$

Since $f(x) = f(-x)$ (even function),

$$\int_{-\infty}^{+\infty} dx e^{\pm i(k-k')x} = 2\pi \delta(k-k')$$

$$\int_{-\infty}^{+\infty} dx e^{\pm ikx} = 2\pi \delta(k)$$

Similarly

$$\int_{\mathbb{R}^3} e^{i\vec{k}\cdot\vec{r}} d^3x = (2\pi)^3 \delta^3(\vec{k})$$

$$\int_{\mathbb{R}^3} e^{i\vec{k}\cdot(\vec{r}-\vec{r}')} d^3k = (2\pi)^3 \delta^3(\vec{r}-\vec{r}')$$

* Suppose $g(x)$ is a function and $g(x_i) = 0$. Then

$$\delta(g(x)) = \sum_i \frac{1}{|g'(x_i)|} \delta(x-x_i)$$

e.g. $g(x) = x^2 - 1$, $x_i = \pm 1$

$$\delta(x^2 - 1) = \sum_i \frac{1}{|g'(x_i)|} \delta(x-x_i) = \frac{1}{2 \cdot (1)} \delta(x-1) + \frac{1}{2 \cdot (-1)} \delta(x+1)$$

$$= \frac{1}{2} (\delta(x-1) - \delta(x+1))$$

$$\therefore \int_{-\infty}^{+\infty} f(x) \delta(x^2 - 1) dx = \frac{1}{2} \int_{-\infty}^{+\infty} f(x) \delta(x-1) dx - \frac{1}{2} \int_{-\infty}^{+\infty} f(x) \delta(x+1) dx$$

$$= \frac{1}{2} (f(1) - f(-1))$$

EX Using $\int_{-\infty}^{+\infty} e^{-ikx} dx = 2\pi \delta(k)$,

$$\int_{-\infty}^{+\infty} x e^{-ikx} dx = \frac{1}{(-i)} \int_{-\infty}^{+\infty} \frac{d}{dk} (e^{-ikx}) dx = i \frac{d}{dk} \int_{-\infty}^{+\infty} e^{-ikx} dx$$

$$= 2\pi i \delta'(k)$$

$$\int_{-\infty}^{+\infty} f(x) \delta'(x) dx = -f'(0)$$

Integration by part