

117. CH.5. Magnetostatics.

▷ Lorentz Force on a moving charge in \vec{B}

$$\vec{F}_q = q (\vec{v} \times \vec{B}).$$

velocity dependent force, unique!

$$[B] = \frac{[F]}{[q][v]} = \frac{N}{C \cdot (m/s)} = \frac{N \cdot s}{C \cdot m} = T \text{ (tesla)}.$$

m $\vec{F}_q = q \vec{E}$: force on a charge in \vec{E}

$$\therefore [B] = \frac{[E]}{[v]} = \frac{(N/C)}{(m/s)} = \frac{N \cdot s}{C \cdot m}.$$

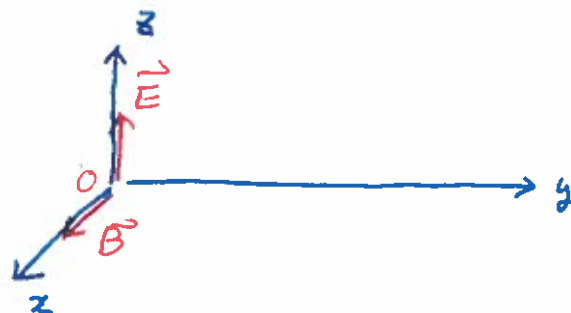
$$\therefore \vec{F}_q = q (\vec{E} + \vec{v} \times \vec{B}).$$

EX Cycloid Motion.

A positive charge is released at O.

$$\vec{v}(0) = 0, \quad \vec{r}(0) = 0.$$

$$\vec{E} = E_0 \hat{z} \quad \text{and} \quad \vec{B} = B_0 \hat{x}.$$



$$\therefore \vec{F}_q = q (\vec{E} + \vec{v} \times \vec{B}).$$

$$\vec{v} \times \vec{B} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ v_x & v_y & v_z \\ B_0 & 0 & 0 \end{vmatrix} = B_0 v_z \hat{y} - B_0 v_y \hat{z}.$$

$$\therefore \vec{F}_q = q (E_0 \hat{z} + B_0 v_z \hat{y} - B_0 v_y \hat{z}) = q B_0 v_z \hat{y} + q (E_0 - v_y B_0) \hat{z}.$$

$$m \ddot{y} = q B_0 \dot{z} \quad \text{and} \quad m \ddot{z} = q (E_0 - B_0 \dot{y}).$$

Define cyclotron frequency $\omega \equiv \frac{q B_0}{m}$ (check if ω has dimension of t^{-1}).

$$\left. \begin{aligned} \ddot{y} &= \omega \dot{z} \\ \ddot{z} &= \omega (E_0 - \dot{y}) \end{aligned} \right\} (*)$$

$$\ddot{y} = \omega \dot{z} \rightarrow \text{into } (*)$$

$$\ddot{y} = -\omega^2 (y - \frac{E_0}{B_0}) \quad \text{--- } \textcircled{a}$$

By substituting $y - \frac{E_0}{B_0} \equiv u$

Eg. \textcircled{a} becomes. $\ddot{u} + \omega^2 u = 0$. then you can find.

$$y(t) = C_1 \cos \omega t + C_2 \sin \omega t + (\frac{E_0}{B_0})t + C_3$$

$$z(t) = C_2 \cos \omega t - C_1 \sin \omega t + C_4$$

Initial conditions $\vec{v}(0) = \vec{r}(0) = 0$.

$$y(0) = C_1 + C_3 = 0 \Rightarrow C_3 = -C_1$$

$$z(0) = C_2 + C_4 = 0 \Rightarrow C_4 = -C_2$$

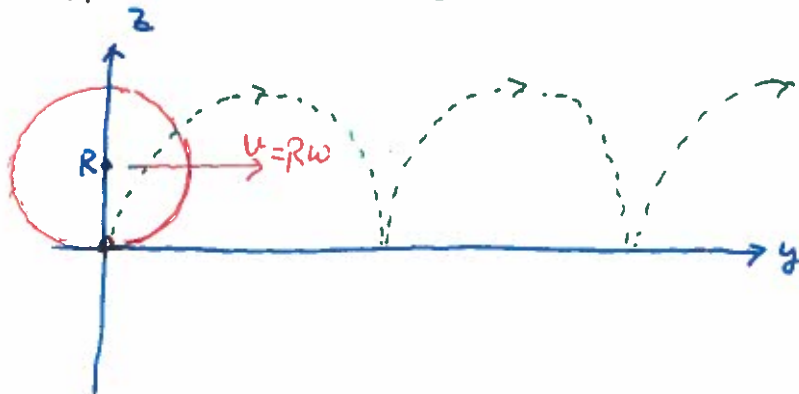
$$\dot{y}(0) = \omega C_2 + \frac{E_0}{B_0} = 0$$

$$\dot{z}(0) = -\omega C_1 = 0$$

$$\left. \begin{array}{l} C_1 = C_3 = 0 \\ C_2 = -\frac{E_0}{\omega B_0} \\ C_4 = \frac{E_0}{\omega B_0} \end{array} \right\}$$

$$\therefore \left. \begin{array}{l} y(t) = \frac{E_0}{\omega B_0} (\omega t - \sin \omega t) \\ z(t) = \frac{E_0}{\omega B_0} (1 - \cos \omega t) \end{array} \right\} \text{ cycloid.}$$

$$(y - R\omega t)^2 + (z - R)^2 = R^2 \quad \text{where } R = \frac{E_0}{\omega B_0}.$$



- Work done by Lorentz force.
charge q moves by $d\vec{l} = \vec{v} dt$ in \vec{B} .

$$\vec{F}_q = q (\vec{v} \times \vec{B})$$

$$\therefore dW = \vec{F}_q \cdot d\vec{l} = q (\vec{v} \times \vec{B}) \cdot \vec{v} dt = 0$$

↑
mutually orthogonal

Magnetic forces do no work.

HW 5.1 & 5.3

- Lorentz force on a current-carrying wire.

Current $I = \frac{dq}{dt}$ and its direction is \vec{B} in the direction of (+) charge
 { in the opposite direction of (-) charge.



If a line charge λ travels along a wire at speed v ,
 then $I = \frac{\lambda v \Delta t}{\Delta t} = \lambda v$. Therefore $\vec{I} = \lambda \vec{v}$
↑
local velocity

Then the force on the wire in \vec{B} ,

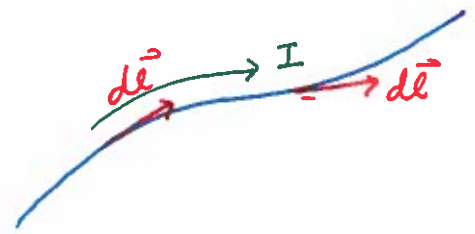
$$\boxed{\vec{F} = \int (\vec{v} \times \vec{B}) dq = \int (\vec{v} \times \vec{B}) \lambda dl = \lambda \int (\vec{v} \times \vec{B}) dl}$$

$$\boxed{= \int (\vec{I} \times \vec{B}) dl}$$

- (i) For a thin wire

$$d\vec{l} \parallel \vec{I}$$

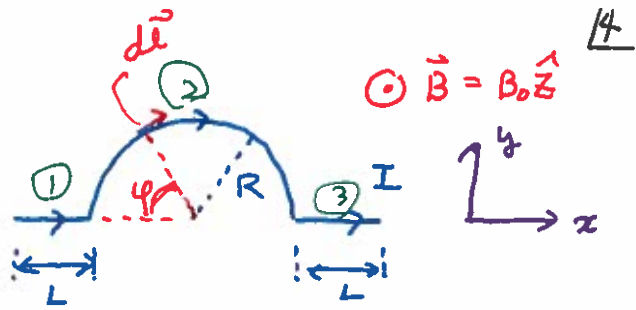
$$\therefore \boxed{\vec{F}_m = \int I (d\vec{l} \times \vec{B})}$$



Ex

Three sections of the wire

- ① & ③ straight
- ② half circle.



$$\vec{F}_1 = \vec{F}_2 = (\vec{I} \times \vec{B})L = I(\hat{i} \times \hat{z})BL = -IBL\hat{j}$$

$$\vec{F}_0 = I \int (\vec{dl} \times \vec{B}) = I \int_0^\pi R d\varphi B_0 (\hat{\varphi} \times \hat{z}) = -I \int_0^\pi R B_0 d\varphi (-\hat{s})$$

Cylindrical coordinates (s, \varphi, z)

$$\hat{s} = -\cos\varphi \hat{x} + \sin\varphi \hat{y}$$
$$= -IRB_0 \int_0^\pi (-\cos\varphi \hat{x} + \sin\varphi \hat{y}) d\varphi$$
$$= -2IRB_0 \hat{y}$$

You already know this by symm.

$$\therefore \vec{F}_m = -IB_0(2R+2L)\hat{y}$$

Simply the total straight length between the end points.



$$|\vec{F}_m| = IB_0 L$$

(ii) Surface current with \vec{K} (surface current density).

$$\vec{K} = \sigma \vec{v} = \frac{d\vec{I}}{dl_\perp} \quad (\text{current per unit width}).$$

$$\therefore \vec{F}_m = \int (\vec{K} \times \vec{B}) \underbrace{da}_{dl dl_\perp} = \int (\vec{v} \times \vec{B}) \sigma da$$



(ii) Volume Current with \vec{J} (vol. current density).

$$\vec{J} = \frac{d\vec{I}}{dA} \quad (\text{current per unit area})$$

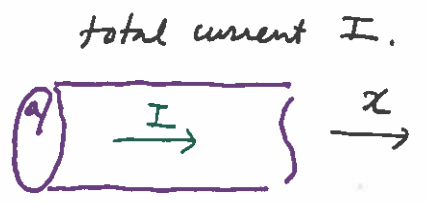
$$= \rho \vec{v}$$

$$\therefore \vec{F}_m = \int (\vec{J} \times \vec{B}) d\tau = \int \rho (\vec{v} \times \vec{B}) d\tau$$

EX.

(a) Uniform current density

$$\vec{J} = \frac{I}{\pi a^2} \hat{x}$$



(b) If $J = kS$, then

$$I = \int_0^a s ds \int_0^{2\pi} d\phi (kS)$$

$$= 2\pi k \left(\frac{a^3}{3}\right)$$



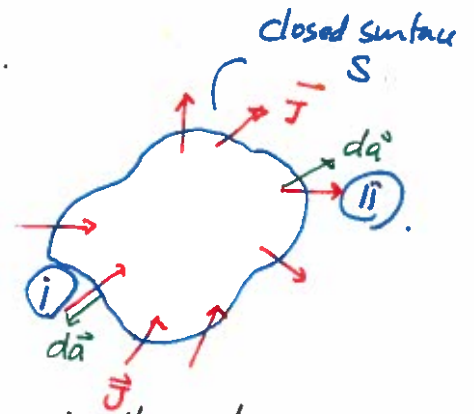
$$\therefore k = \frac{3I}{2\pi a^3}$$

Continuity Eq. ("Charge Conservation")

$$I = \int_S \vec{J} \cdot d\vec{a} = \int_S \vec{J} \cdot \vec{n} \, da \quad (\text{flux})$$

For a closed surface,

$$\oint_S \vec{J} \cdot \vec{a} = \int_V (\vec{\nabla} \cdot \vec{J}) d\tau$$



But this is related to charge accumulation in the volume.

- (i) $\vec{J} \cdot \vec{a} < 0$ will increase the charge in the volume.
- (ii) $\vec{J} \cdot \vec{a} > 0$ will decrease " " " "

$$\therefore \int_V (\nabla \cdot \vec{J}) d\tau = -\frac{dQ_{en}}{dt} = -\frac{d}{dt} \int_V \rho d\tau = -\int \frac{\partial \rho}{\partial t} d\tau.$$

therefore $\boxed{\frac{\partial \rho}{\partial t} + \nabla \cdot \vec{J} = 0}$ Continuity Eq. ($\vec{J} = \rho \vec{v}$).

Any volume density ρ_K ($K = \text{mass, charge, } \dots$).

Then the associated current density $\vec{J}_K = \rho_K \vec{v}$ local velocity.

If $\rho_K V = K$ (mass, charge, ...) is conserved,

$$\frac{\partial \rho_K}{\partial t} + \nabla \cdot (\rho_K \vec{v}) = \frac{\partial \rho_K}{\partial t} + \nabla \cdot \vec{J}_K = 0.$$

HW 5.4, 3.5, 5.6.