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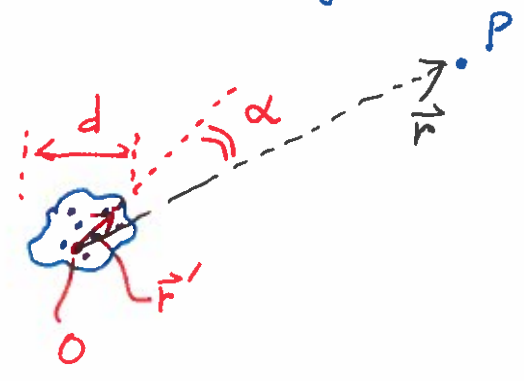
► Multipole Expansion

$V(\vec{r})$  or  $\vec{E}(\vec{r})$  from a localized charge distribution.  
for  $r \gg d$  where  $d$  is typical size of the localized charge.

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \rho(\vec{r}') \cdot \frac{1}{r} d\tau'$$

$$r = |\vec{r} - \vec{r}'|$$

$$Q = \int \rho(\vec{r}') d\tau'$$



For  $r \gg d$ ,

$$V(\vec{r}) \approx \frac{1}{4\pi\epsilon_0} \cdot \frac{Q}{r} \quad \text{behaves like a pt. charge for } r \gg d$$

However, if  $Q = 0$  but  $\rho(\vec{r}') \neq 0$ , then

$$V(\vec{r}) \sim \frac{1}{r^n} \quad (n \geq 2), \quad \text{or } E \sim \frac{1}{r^{n+1}} \quad (n \geq 2)$$

+	+ -	+ - - +	
•	• •	• •	• •
$Q = +$	$Q = 0$	$Q = 0$	$Q = 0$
$V \sim \frac{1}{r}$	$\frac{1}{r^2}$	$\frac{1}{r^3}$	$\frac{1}{r^4}$
$E \sim \frac{1}{r^2}$	$\frac{1}{r^3}$	$\frac{1}{r^4}$	$\frac{1}{r^5}$
<b>Monopole</b>	<b>Dipole</b>	<b>Quadrupole</b>	<b>Octopole</b>

||| leading term

Taylor Expansion!  $r \gg d$ .

$$r^2 = r^2 + (r')^2 - 2rr' \cos\alpha \quad \cos\alpha = \angle(\vec{r}', \vec{r})$$

Pay attention to the location of the origin!

$$r' \leq d \quad \therefore r \gg r'$$

$$r^2 = r^2 \left[ 1 + \left(\frac{r'}{r}\right)^2 - 2\left(\frac{r'}{r}\right)\cos\alpha \right] = r^2(1 + \epsilon) \quad \underline{12}$$

where  $\epsilon = \left(\frac{r'}{r}\right) \left(\frac{r'}{r} - 2\cos\alpha\right)$

$$\therefore r = r \sqrt{1 + \epsilon} \quad \text{and}$$

$$\begin{aligned} \frac{1}{r} &= \frac{1}{r} \cdot \frac{1}{\sqrt{1 + \epsilon}} = \frac{1}{r} \left( 1 - \frac{1}{2}\epsilon + \frac{3}{8}\epsilon^2 - \frac{5}{16}\epsilon^3 + \dots \right) \\ &= \frac{1}{r} \left[ 1 - \frac{1}{2}\left(\frac{r'}{r}\right)\left(\frac{r'}{r} - 2\cos\alpha\right) + \frac{3}{8}\left(\frac{r'}{r}\right)^2\left(\frac{r'}{r} - 2\cos\alpha\right)^2 - \frac{5}{16}\left(\frac{r'}{r}\right)^3\left(\frac{r'}{r} - 2\cos\alpha\right)^3 + \dots \right] \\ &= \frac{1}{r} \left[ 1 + \underbrace{\left(\frac{r'}{r}\right)\cos\alpha - \frac{1}{2}\left(\frac{r'}{r}\right)^2}_{\substack{\cancel{2} \\ \cancel{2}}} + \underbrace{\frac{3}{8} \cdot 4\cos^2\alpha \left(\frac{r'}{r}\right)^2 + \frac{3}{8} \cdot (-4) \left(\frac{r'}{r}\right)^3 \cos\alpha + \frac{3}{8} \left(\frac{r'}{r}\right)^4}_{\substack{\cancel{2} \\ \cancel{2}}} \right. \\ &\quad \left. - \frac{5}{16} \left(\frac{r'}{r}\right)^3 (-8)\cos^3\alpha - \frac{5}{16} (3 \cdot 2^2) \left(\frac{r'}{r}\right)^4 \cos^2\alpha - \frac{5}{16} (-2) \cdot 3 \left(\frac{r'}{r}\right)^5 \cos\alpha - \frac{5}{16} \left(\frac{r'}{r}\right)^6 + \dots \right] \end{aligned}$$

$$= \frac{1}{r} \left[ \underbrace{1}_{P_0(\cos\alpha)} + \underbrace{\left(\frac{r'}{r}\right)\cos\alpha}_{P_1(\cos\alpha)} + \underbrace{\left(\frac{3\cos^2\alpha - 1}{2}\right)\left(\frac{r'}{r}\right)^2}_{P_2(\cos\alpha)} + \underbrace{\left(\frac{5\cos^3\alpha - 3\cos\alpha}{2}\right)\left(\frac{r'}{r}\right)^3}_{P_3(\cos\alpha)} + \dots \right]$$

$$= \frac{1}{r} \left[ P_0(\cos\alpha) + \left(\frac{r'}{r}\right) P_1(\cos\alpha) + \left(\frac{r'}{r}\right)^2 P_2(\cos\alpha) + \left(\frac{r'}{r}\right)^3 P_3(\cos\alpha) + \dots \right]$$

$$= \frac{1}{r} \sum_{n=0}^{\infty} \left(\frac{r'}{r}\right)^n P_n(\cos\alpha)$$

Multipole Expansion of Legendre Polynomials

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{r^n}{r^{n+1}} \int (r')^n P_n(\cos\alpha) \rho(\vec{r}') d\tau'$$

$$= \frac{1}{4\pi\epsilon_0} \left[ \underbrace{\frac{1}{r} \int \rho(\vec{r}') d\tau'}_{\text{monopole}} + \underbrace{\frac{1}{r^2} \int r' \rho(\vec{r}') P_1(\cos\alpha) d\tau'}_{\text{dipole}} \right. \\ \left. + \frac{1}{r^3} \int (r')^2 \rho(\vec{r}') P_2(\cos\alpha) d\tau' \right. \\ \left. + \dots \right]$$

← quadrupole

EX Prob 3.28

Uniform ring charge with  $\lambda$  of radius  $R$

$$V(r) = \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{1}{r^{n+1}} \int (r')^n P_n(\cos\alpha) \rho(r') d\tau'$$

$$\vec{r} = (r, \theta, \varphi) = r \sin\theta \cos\varphi \hat{x} + r \sin\theta \sin\varphi \hat{y} + r \cos\theta \hat{z}$$

$$\vec{r}' = (R, \frac{\pi}{2}, \varphi') = R \cos\varphi' \hat{x} + R \sin\varphi' \hat{y} + R \cos\theta \hat{z}$$

$$\vec{r} \cdot \vec{r}' = r r' \cos\alpha = r R (\underbrace{\sin\theta \cos\varphi \cos\varphi' + \sin\theta \sin\varphi \sin\varphi'}_{\cos\alpha})$$

$$\int_V \rho(r') d\tau' \rightarrow \int_C \lambda(r') dl' = \int_0^{2\pi} \lambda R d\varphi'$$

For  $n=0$  (monopole)

$$V_0 = \frac{1}{4\pi\epsilon_0} \frac{1}{r} \int_0^{2\pi} (1) \lambda R d\varphi' = \frac{\lambda R}{2\epsilon_0} \left(\frac{1}{r}\right)$$

For  $n=1$  (dipole)

$$V_1 = \frac{1}{4\pi\epsilon_0} \frac{1}{r^2} \int_0^{2\pi} (r')^1 P_1(\cos\alpha) \lambda R d\varphi'$$

$$= \frac{1}{4\pi\epsilon_0} \left(\frac{1}{r^2}\right) \lambda R^2 \int_0^{2\pi} (\sin\theta \cos\varphi \cos\varphi' + \sin\theta \sin\varphi \sin\varphi') d\varphi'$$

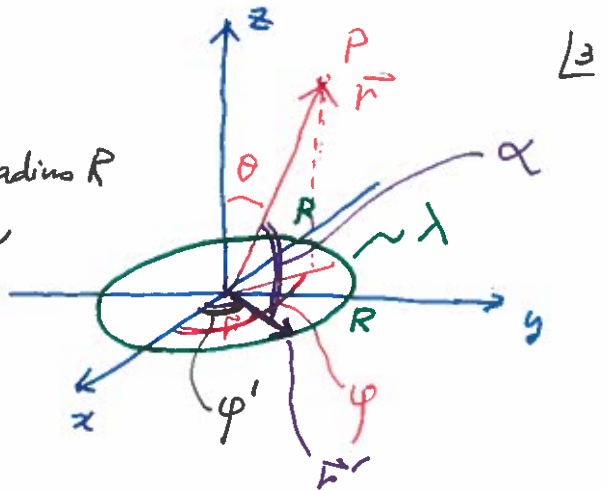
= 0 (no dipole)

For  $n=2$

$$V_2 = \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} \int_0^{2\pi} (r')^2 P_2(\cos\alpha) \lambda R d\varphi'$$

$$= \frac{1}{4\pi\epsilon_0} \frac{\lambda R^3}{r^3} \int_0^{2\pi} \frac{1}{2} (3\cos^2\alpha - 1) d\varphi' = \frac{1}{4\pi\epsilon_0} \left(\frac{\lambda R^3}{2r^3}\right) \int_0^{2\pi} \{3\sin^2\theta (\cos\varphi \cos\varphi' + \sin\varphi \sin\varphi')^2 - 1\} d\varphi'$$

$$* \int_0^{2\pi} \sin^2\theta d\theta = \int_0^{2\pi} \cos^2\theta d\theta = \pi \quad \text{and} \quad \int_0^{2\pi} \sin\theta \cos\theta d\theta = 0.$$



$$\begin{aligned} \therefore V_2 &= \frac{1}{4\pi\epsilon_0} \left( \frac{\lambda R^3}{2r^3} \right) \left[ 3 \sin^2\theta \left( \underbrace{\pi \cos^2\varphi + \pi \sin^2\varphi}_{\pi} \right) - 2\pi \right] \\ &= \frac{1}{4\pi\epsilon_0} \left( \frac{\lambda R^3}{2r^3} \right) \pi (3 \sin^2\theta - 2) \\ &= \frac{1}{8\epsilon_0} \left( \frac{\lambda R^3}{r^3} \right) (3 \sin^2\theta - 2) \\ &= -\frac{\lambda}{8\epsilon_0} \left( \frac{R}{r} \right)^3 (3 \cos^2\theta - 1) = -\frac{\lambda}{4\epsilon_0} \left( \frac{R}{r} \right)^3 P_2(\cos\theta) \end{aligned}$$

**HW** 3.27

▷ Dipole moment.

$\vec{p} \equiv \int \vec{r}' \rho(\vec{r}') d\tau'$  for continuous charge distribution

then  $V_{\text{dipole}} = \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} \int r' \rho(\vec{r}') P_1(\cos\alpha) d\tau'$   
 $= \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \hat{r}}{r^2}$

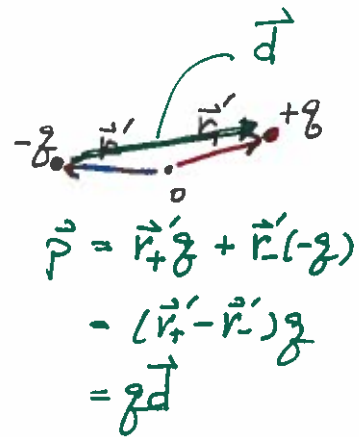
For a collection of pt charges  $q_i$ 's at  $\vec{r}_i$ '

$$\vec{p} = \sum_{i=1}^n \vec{r}_i' q_i$$

$$|\vec{p}| \approx qd$$

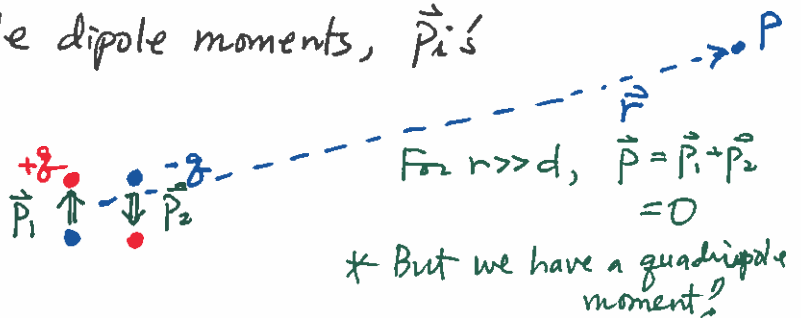
Perfect dipole is a mathematical object  $d \rightarrow 0$  while  $qd = p \neq 0$ .

In reality, if  $|d| \ll r$ ,  $\vec{p} = q\vec{d}$  can work as a perfect dipole!



One can have multiple dipole moments,  $\vec{p}_i$ 's

$$\therefore \vec{p} = \sum_{i=1}^n \vec{p}_i$$



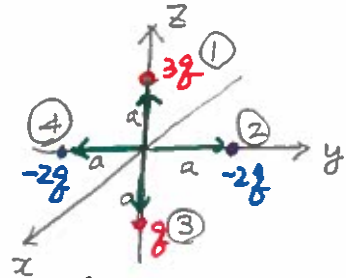
EX What is  $V(\vec{r})$  for  $r \gg a$

$$Q = \sum_i q_i = 0 \quad (\text{no monopole}).$$

$$\vec{p} = \sum_{i=1}^4 \vec{r}'_i q_i$$

$$= \{3q(+\hat{z}) + (-2q)(+\hat{y}) + q(-\hat{z}) - 2q(-\hat{y})\} a$$

$$= 2qa\hat{z}$$



**HW** 3.30

⊙ Dipole field  $\vec{E}_{\text{dip}}$

$$V_{\text{dip}}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \hat{r}}{r^2} \quad (\text{see the geometry})$$

$$(\vec{p} = p\hat{z}) = \frac{p}{4\pi\epsilon_0} \frac{\cos\theta}{r^2}$$

$$\therefore \vec{E}_{\text{dip}} = -\vec{\nabla} V_{\text{dip}}$$

$$= -\frac{\partial V}{\partial r} \hat{r} - \frac{1}{r} \frac{\partial V}{\partial \theta} \hat{\theta} - \frac{1}{r \sin\theta} \frac{\partial V}{\partial \varphi} \hat{\varphi}$$

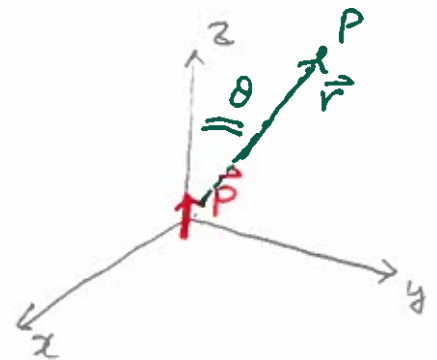
$$= \frac{2p \cos\theta}{4\pi\epsilon_0} \left(\frac{1}{r^3}\right) \hat{r} + \frac{p \sin\theta}{4\pi\epsilon_0} \left(\frac{1}{r^3}\right) \hat{\theta} + 0$$

$$\vec{E}_{\text{dip}}(r, \theta) = \frac{p}{4\pi\epsilon_0 r^3} (2\cos\theta \hat{r} + \sin\theta \hat{\theta})$$

$$\text{for } \theta=0 \quad \vec{E}_{\text{dip}} = \frac{p}{2\pi\epsilon_0} \left(\frac{1}{r^3}\right) \hat{r}$$

$$\text{for } \theta=\frac{\pi}{2} \quad \vec{E}_{\text{dip}} = \frac{p}{4\pi\epsilon_0} \left(\frac{1}{r^3}\right) \hat{\theta}$$

azimuthal symm.  
no  $\varphi$ -dep.



**HW** 3.35, 3.36, 3.52 if you want to (a) (b) (c) only